

Homework 3

Mathematical Methods I: Fall 2017

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1. Problem 1 - Elastic Rods:

- (a) Show that The elastic energy in a bent beam is is

$$U[y] = \int_0^L \frac{1}{2} YI (y'')^2 dz$$

given that the elastic energy per unit length of a bent steel rod is given by $\frac{1}{2} \frac{YI}{R^2}$ where R is the radius of curvature due to bending.

Proof. From vector calculus we know that the length of the radius of curvature vector for a curve $y(x)$ is given by

$$|\vec{R}| = \frac{(1 + y'^2)^{\frac{3}{2}}}{y''} \approx \frac{1}{y''}$$

where we approximate y'^2 to be small enough to ignore in this case. Then we have

$$U[y] = \int_0^L \frac{1}{2} \frac{YI}{R^2} dz = \int_0^L \frac{1}{2} \frac{YI}{(y'')^2} dz$$

□

- (b) Show that if there is a load of mass M on top of the rod, the energy can be approximated by

$$U[y] = \int_0^L \left(\frac{YI}{2} (y'')^2 - \frac{Mg}{2} (y')^2 \right) dz$$

Proof. Gravitational potential energy is clearly going to be

$$U_g = MgL_z$$

where L_z is the height of the load after the rod bends. We can calculate this by:

$$L_z = \int dz$$

along the curve of the rod. We know that $dl = \sqrt{1 + y'^2} dz$, so we plug this into

$$\int dz = \int \frac{dl}{\sqrt{1 + y'^2}}$$

But now we assume that the deflection is very small so we get that $dl \approx dz$ and expand the denominator:

$$U_g = \int_0^L Mg dz \approx \int_0^L Mg \left(1 - \frac{1}{2} y'^2 \right) dz$$

Thus the total energy functional is of the desired form (excluding the constant term MgL). □

- (c) Show that the column is unstable to buckling and collapses when $Mg \geq \frac{\pi^2}{L^2} YI$.

Proof. Plugging in the ansatz for the solutions:

$$y(z) = \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi z}{L}\right)$$

into the energy functional:

$$U = \sum_{n=1}^{\infty} a_n^2 \int_0^L \frac{YI}{2} \left(\frac{n\pi}{L}\right)^4 \sin^2\left(\frac{n\pi z}{L}\right) - \frac{Mg}{2} \left(\frac{n\pi}{L}\right)^2 \cos^2\left(\frac{n\pi z}{L}\right) dz$$

Doing the trig integrals out, we get factors of $\frac{1}{2}$ out, and the remaining terms depend on n . We want to know when these terms are negative for a given n value, which would tell us that the energy drops down to a negative value. This first happens at $n = 1$ as the coefficient becomes negative when

$$Mg \geq YI \left(\frac{n\pi}{L}\right)^2$$

□

- (d) The light cantilever: Find $y(z)$ for $0 < z < L$ assuming that a rod is fixed into a wall with a load of mass M hanging at the end.

Proof. We want to minimize the energy functional

$$U = \int_0^L \left(\frac{YI}{2} (y'')^2\right) dz + Mgy(L)$$

But we will be careful not to throw out terms when we integrate by parts.

$$\begin{aligned} U(y + \delta y) - U(y) &= \int_0^L \frac{YI}{2} \left((y + \delta y)^2 - (y'')^2 \right) dz + Mg(y(L) + \delta y(L) - y(L)) \\ \delta U &= \int_0^L \frac{YI}{2} (2y''(\delta y)'' + (\delta y'')^2) dz + Mg\delta y(L) \\ \delta U &= \int_0^L \frac{YI}{2} (2y''(\delta y)'' + \mathcal{O}(\delta y)^2) dz + Mg\delta y(L) \\ \delta U &= \int_0^L YI(y''(\delta y)'' dz + Mg\delta y(L) \end{aligned}$$

We now integrate by parts twice:

$$\begin{aligned} \delta U &= (y''\delta y)|_0^L - \int_0^L y^{(3)}(\delta y)' + Mg\delta y(L) \\ \delta U &= (y''\delta y)|_0^L - (y^3\delta y)|_0^L + \int_0^L y^{(4)}\delta y dz + Mg\delta y(L) = 0 \end{aligned}$$

This is true for any $\delta y(L)$, knowing that $\delta y(0) = 0$. If we set all these terms to zero, and factor out the terms dependent on we the differential equation

$$y^{(4)} = 0$$

with the boundary conditions

$$\begin{aligned} Mg &= YIy^{(3)} \\ y''(L) &= 0 \\ y'(0) &= y(0) = 0 \end{aligned}$$

The most general solution to the differential equation is $y(z) = Az^3 + Bz^2 + Cz + D$ But we know immediately that D is zero. From the other boundary conditions we get:

$$y''(L) = 6AL + 2B = 0$$

$$6A = \frac{Mg}{YI}$$

Thus the solution is

$$y(z) = \frac{Mg}{YI} \left(\frac{1}{6}z^3 - \frac{L}{2}z^2 \right)$$

And

$$y(L) = -\frac{MgL^2}{3YI}$$

□

2. Lagrange Multipliers

- (a) Find the stationary points of the function

$$f(x, y) = 13x^2 + 8xy + 7y^2$$

subject to $x^2 + y^2 = 1$.

Proof. We first express this as a matrix multiplication:

$$\mathbf{x}^t \mathbf{A} \mathbf{x} = f(x, y)$$

Minimizing this function with the constraint gives us an eigenvalue problem to solve:

$$\langle \mathbf{x}, \mathbf{A} \mathbf{x} \rangle - \lambda (\langle \mathbf{x}, \mathbf{x} \rangle - 1) = g(\mathbf{x})$$

We then differentiate:

$$\frac{\partial g}{\partial \mathbf{x}} = 2\mathbf{A} \mathbf{x} - 2\lambda \mathbf{x} = 0$$

And arrive at

$$\mathbf{A} \mathbf{x} = \lambda \mathbf{x}$$

. Then we can find the eigenvalues of this matrix using the standard techniques. We find them to be

$$\lambda = 10, 5$$

We can also find normalized eigenvectors and we find them to be

$$e_1 = \frac{1}{\sqrt{5}}(2, 1)$$

and

$$e_2 = \frac{1}{\sqrt{5}}(-1, 2)$$

This gives us our two stationary points. But since we know the constraint is in terms of x^2 and y^2 , we actually get 4 stationary points that are

$$(x, y) = \pm e_1, \pm e_2$$

□

3. The Catenary Again:

- (a) From the resulting functional derivative, derive two coupled equations for the catenary, one for $x(s)$ and one for $y(s)$.

Proof. We have to minimize the energy functional

$$U(x, y) = \int_0^L \rho g y(s) ds + \int_0^L (\dot{x}^2 + \dot{y}^2 - 1) \lambda(s) ds$$

where s is the coordinate along the curve. We can use euler-lagrange equations

$$\begin{aligned} \frac{\partial f}{\partial x} &= \frac{d}{ds} \left(\frac{\partial f}{\partial \dot{x}} \right) \\ \frac{\partial f}{\partial y} &= \frac{d}{ds} \left(\frac{\partial f}{\partial \dot{y}} \right) \Rightarrow \\ 0 &= \dot{\lambda} \dot{x} + \lambda \ddot{x} \\ \rho g &= 2 \dot{\lambda} \dot{y} + 2 \lambda \ddot{y} \end{aligned}$$

Now introduce $\dot{x} = \cos \psi$, $\dot{y} = \sin \psi$

$$\begin{aligned} 0 &= \dot{\lambda} \cos \psi - \lambda \sin \psi \dot{\psi} \\ \rho g &= 2(\dot{\lambda} \sin \psi + \lambda \cos \psi \dot{\psi}) \end{aligned}$$

Square these and add them up, we get:

$$(\rho g)^2 = 4(\dot{\lambda}^2 + \lambda^2 \dot{\psi}^2)$$

From looking at a section of chain, we can deduce that

$$T(s + ds)_y - T(s)_y = \rho g ds$$

and

$$T(s)_x = T(s + ds)_x$$

If we expand $T(s + ds)_y = T(s)_y + \dot{T}_y ds$ and plug in, we find that $\dot{T}_y = \rho g$. We can then define $T_x = 2\lambda \cos \psi$ and $T_y = 2\lambda \sin \psi$ such that $\dot{T}_y^2 + \dot{T}_x^2 = (\rho g)^2$. \square

- (b) Now find the material density $\rho(s)$ in order for a length of chain $\frac{\pi a}{2}$ to hang in an arc of a circle of radius a .

Proof. If we draw the arc for ψ along the arc of the circle, we can deduce that $\psi = \frac{s}{a}$. We also know that $\dot{T}_x = 0$, thus $\frac{d}{ds}(\lambda \cos \psi) = 0$, $\lambda(s) \cos \psi = K$. Thus

$$\rho(s)g = 2\partial_s(\lambda \sin \psi(s)) = 2\partial_s(K \tan(\frac{s}{a})) = \frac{2K}{a} \sec^2(\frac{s}{a})$$

\square